



# Breaking down chipping and fragmentation in sediment transport: the control of material strength

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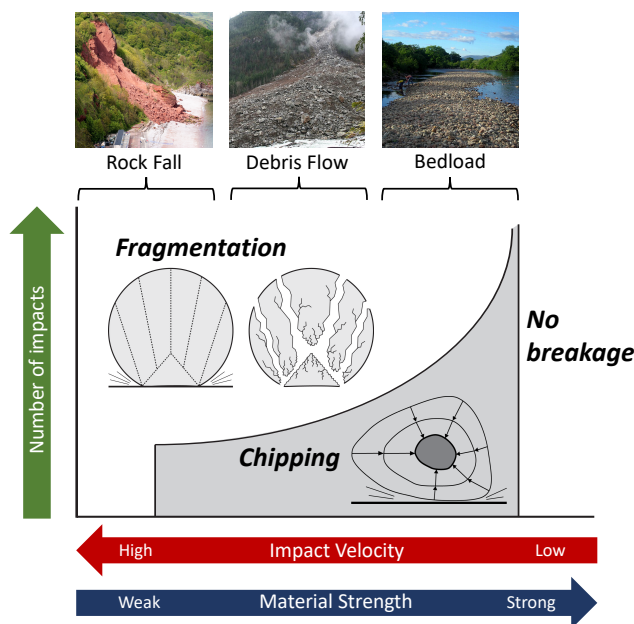
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**Abstract.** As rocks are transported, they primarily undergo two breakdown mechanisms: chipping and fragmentation. Chipping occurs at relatively low collision energies typical of bed-load transport, and involves shallow cracking; this process rounds river pebbles in a universal manner. Fragmentation involves catastrophic breakup by fracture growth in the bulk — a response that occurs at high collision energies such as rock falls — and produces angular shards. Despite its geophysical significance, the transition from chipping to fragmentation is not well studied. Indeed, most models implicitly assume that impact erosion of pebbles and bedrock is governed by fragmentation rather than chipping. Here we experimentally delineate the boundary between chipping and fragmentation by examining the mass and shape evolution of concrete particles in a rotating drum. Attrition rate should be a function of both impact energy and material strength; here we keep the former constant, while systematically varying the latter. For sufficiently strong particles, chipping occurred and was characterized by the following: daughter products were significantly smaller than the parent; attrition rate was independent of material strength; and particles experienced monotonic rounding toward a spherical shape. As strength decreased, fragmentation became more significant: mass of daughter products became larger and more varied; attrition rate was inversely proportional to material strength; and shape evolution fluctuated and became non monotonic. Our results validate a previously proposed probabilistic model for impact attrition, and indicate that bedrock erosion models predicated on fragmentation failure need to be revisited. We suggest that the shape of natural pebbles may be utilized to deduce the breakdown mechanism, and infer past transport environments.

## 1 Introduction

The energy associated with transporting sediment influences the shape of the particles being moved. Events such as rock falls, landslides, or debris flows tend to be highly energetic and cause rocks to fragment into large, angular pieces (Bernd et al., 2010; Arabia and Sklar, 2016). Bed-load transport is lower energy, and causes river rocks to round into smooth pebbles (Krumbein, 1941; Kuenen, 1956; Attal and Lavé, 2009; Domokos et al., 2014). The attrition mechanism that occurs during transport depends on impact energy and material properties, including size, shape, and material strength (Yashima et al., 1987; Tavares and King, 1998; Sklar and Dietrich, 2004; Tavares and das Neves, 2008; Domokos et al., 2015); these attrition mechanisms can be classified using the terms *frictional abrasion*, *chipping*, and *fragmentation* (Bemrose and Bridgwater, 1987; Ghadiri



**Figure 1.** Schematic phase diagram delineating the attrition mechanism as a function of impact velocity or material strength, and number of impacts; adapted from the proposed diagram of Zhang and Ghadiri (2002). For extremely weak materials or high impact velocities (left), explosive fragmentation can occur from a single collision. As material strength increases or impact velocity decreases (center), fragmentation occurs through fatigue failure, where cracks grow progressively through the material with each collision. In the limit of pure chipping, which applies to low impact velocities or strong materials (right), the number of impacts required to produce fragmentation increases toward infinity. Examples of transport mechanisms associated with fragmentation, chipping and fragmentation, and pure chipping are a rock fall at Oddicombe Beach, England, debris flow at Goodell Creek, Washington, USA, and bedload in the River Nevis, Scotland, respectively. Fragmentation figure adapted from Salman et al. (2004), chipping figure adapted from Szabó et al. (2013), transport mechanism images reproduced from Wikimedia Commons.

and Zhang, 2002). This study examines the chipping and fragmentation regimes, and how material strength influences attrition mechanism (Fig. 1).

The characteristic smooth, round shape of pebbles is the result of chipping, when low-energy collisions cause small particles to break off the parent rock (Rogers and Schubert, 1963; Ghadiri and Zhang, 2002; Jerolmack and Brzinski, 2010). Lateral cracking is known to produce chipping in natural rock materials (Wilshaw, 1971). A recent experimental study on binary collisions of rocks (Miller and Jerolmack, 2020) has shown that chipping involves fragmentation within a shallow region, that is bounded from below by surface-parallel cracks (Fig. 2). These locally shattered regions are consistent with Hertzian fracture cones, and the related phenomenon of conchoidal fracturing, that are typically observed in very brittle glasses and ceramics (Wilshaw, 1971; Greeley and Iversen, 1987; Mohajerani and Spelt, 2010; Wang et al., 2017); though, some plastic deformation may occur also (Rhee et al., 2001; Ghadiri and Zhang, 2002; Wilshaw, 1971). Hertzian fracture has also been proposed as the mechanism of erosion of bedrock by aeolian transport (Greeley and Iversen, 1987). In the chipping regime



of impact attrition, each collision produces a shower of fine particles — meaning that there is a localized near-surface region where yield is exceeded.

Unlike chipping, catastrophic fragmentation occurs when high-energy collisions cause cracks to propagate radially into the bulk. These radial cracks can split the parent rock into irregularly-shaped daughter particles of significant size in addition to smaller fragments (Perfect, 1997; Kun and Herrmann, 1999; Salman et al., 2004; Grady, 2010). At intermediate impact energies, fatigue failure occurs as fractures grow into the rock as a result of repeated impacts (Bitter, 1963; Moss et al., 1973). We note that this fatigue failure mechanism forms the basis for widely used bedrock erosion models (Sklar and Dietrich, 2001, 2004).

Since breakdown mechanism and the resulting particle shape vary depending on transport energy and material strength, particle shape can be used to infer transport history. Previous studies on attrition by chipping have shown that pebble shape can be used to determine transport distance in fluvial environments (e.g., Attal and Lavé, 2009; Szabó et al., 2013, 2015). Similar studies have examined attrition mechanisms in a dune field (e.g., Jerolmack et al., 2011) and on a clastic beach (e.g. Bertoni et al., 2016). However, fewer studies have investigated particle shape change in high energy environments (e.g., Bernd et al., 2010; Arabnia and Sklar, 2016) or examined the transition between chipping and fragmentation at intermediate energies (e.g., Moss et al., 1973; Adams, 1979). Novák-Szabó et al. (2018) proposed that there is a critical energy associated with this transition, and that natural rock materials in bed load are far below this value. This study utilizes laboratory experiments to examine chipping and fragmentation in order to better understand the connection between breakdown mechanism and shape evolution. As a result, particle shape can be more accurately used to determine past transport conditions and environments.

### 1.1 Attrition and the Shape of River Rocks

As sediment is produced in upland regions and makes its way downstream in rivers, particles transform from blocky, angular clasts to smooth, ellipsoidal pebbles (Krumbein, 1941; Kuenen, 1956; Parker, 1991; Kodama, 1994; Attal and Lavé, 2009; Domokos et al., 2014). This shape change is associated with attrition, and is expected to result in an exponentially-decreasing pebble volume,  $V$ , with distance downstream,  $x$  (Kodama, 1994; Lewin and Brewer, 2002), as expressed in the volumetric version (Novák-Szabó et al., 2018) of Sternberg's Law (Sternberg, 1875):

$$V = V_o e^{-\gamma x}, \quad (1)$$

where  $V_o$  is initial particle volume at the upstream boundary, and  $\gamma$  is a diminution coefficient. If one assumes that particle diameter  $D \propto V^{\frac{1}{3}}$  — an assumption that is likely violated for attrition of angular particles (Domokos et al., 2014) — Sternberg's Law can also be cast in terms of particle diameter (Sternberg, 1875; Kodama, 1994; Lewin and Brewer, 2002; Attal and Lavé, 2009). Unfortunately, size-selective sorting is also expected to produce exponential downstream fining (Paola et al., 1992; Seal and Paola, 1995; Paola and Seal, 1995; Ferguson et al., 1996; Gasparini et al., 1999; Fedele and Paola, 2007). The relative importance of attrition versus sorting as the dominant downstream fining mechanism has been debated, and can vary depending on transport conditions, sediment supply, lithology, and other environmental factors (e.g., Parker, 1991; Kodama, 1992; Mikoš, 1994; Lewin and Brewer, 2002; Szabó et al., 2013; Miller et al., 2014). Recent studies, however, have shown how downstream



changes in pebble shape alone may be used to isolate the contribution of impact attrition to downstream fining — assuming that all mass loss results from chipping (Szabó et al., 2015; Novák-Szabó et al., 2018).

## 1.2 Mass Loss as a Result of Attrition

It is common to cast the attrition process in terms of the mass lost per collision,  $\Delta m$ . This mass loss is proportional to collision energy,  $\Delta E$ , such that:

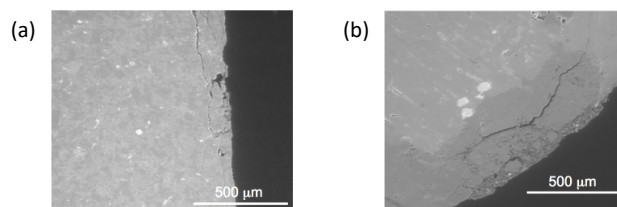
$$\Delta m = C_1 A \Delta E \quad (2)$$

where  $A$  is a material susceptibility parameter (Anderson, 1986) that Miller and Jerolmack (2020) called the “Attrition Number” [ $\text{s}^2 \text{m}^{-2}$ ], and  $C_1$  is an experimentally determined constant. Since collision energy is a function of the mass of the particle,  $m$ , and impact velocity,  $v_i$ , where  $\Delta E = \frac{1}{2} m v_i^2$ , we expect that mass should decrease exponentially with number of impacts (Novák-Szabó et al., 2018). The Attrition Number,  $A$ , incorporates various aspects of material strength that determine susceptibility to attrition, such as hardness, fracture toughness, yield strength, or Young’s modulus (e.g., Ghadiri and Zhang, 2002; Sklar and Dietrich, 2004; Wang et al., 2011).

When the stress produced by a sufficiently powerful impact on a sufficiently weak material exceeds the breaking strength of that material, a crack tip forms. The energy from the surrounding strain can then cause the crack tip to spread through the material and split the rock (Bond, 1952, 1955). If the energy imparted by the collision is insufficient, repeated impacts can cause the crack tips to propagate into the material through fatigue failure (Bitter, 1963; Moss et al., 1973). Fractures formed by brittle failure tend to be large-scale, developing through the volume of the particle as a result of elastic deformation (Evans and Wilshaw, 1977; Ghadiri and Zhang, 2002). It has been proposed that susceptibility to brittle fracture, called the brittle Attrition Number,  $A_b$ , depends on a material’s ability to store energy elastically (Engle, 1978; Wang et al., 2011). From mechanical considerations and dimensional analysis, several studies arrived at a similar parameter (Sklar and Dietrich, 2004; Wang et al., 2011; Miller and Jerolmack, 2020):

$$A_b = \frac{\rho_s Y}{\sigma_s^2}, \quad (3)$$

where  $\sigma_s$  is yield strength [Pa],  $Y$  is Young’s modulus [Pa], and  $\rho_s$  is the density of the material [ $\text{kg m}^{-3}$ ]. While the parameter  $A_b$  formally relates to crack growth in purely brittle materials under fragmentation, it was found to describe mass loss due to shallow chipping of natural (and possibly semi-elastic) rocks (Miller and Jerolmack, 2020). For this reason we will use  $A_b$  to characterize material strength in our study. The applicability of  $A_b$  may imply that chipping of natural rocks is also a form of brittle fracture, but that its depth is limited; the locally shattered regions produced by impacts may be Hertzian fracture cones (Greeley and Iversen, 1987; Wang et al., 2017; Miller and Jerolmack, 2020). Impact experiments on concrete and rocks indicate that this picture is applicable for most materials (Momber, 2004a). It is important to note, however, that plastic deformation may be relevant for impact attrition — even for very brittle materials (Rhee et al., 2001). A ‘brittleness ratio’  $H/K_c$ , where  $H$  is hardness [Pa] and  $K_c$  is fracture toughness [ $\text{N m}^{-3/2}$ ], has been proposed to delineate the transition from purely elastic to semi-elastic deformation (Momber, 2004b). Indeed, for chipping of semi-elastic materials, Ghadiri and Zhang (2002) proposed



**Figure 2.** Scanning electron microscope images depicting subsurface lateral cracks, formed after repeated binary collisions of rocks at energies representative of bed-load transport. We speculate that these cracks, and the shattered region bounded by them, are the result of Hertzian fracture cones formed by impulsive collisions. Images show thin sections of (a) quartz diorite, and (b) volcanoclastic rock. Figure reproduced from Miller and Jerolmack (2020).

an alternative Attrition Number that depends on  $H$  and  $K_c$ , rather than  $Y$  and  $\sigma_s$ . All of these parameters may be correlated with each other, depending on the deformation mechanism (Shipway and Hutchings, 1993; Rhee et al., 2001; Emmerich, 2007; Mohajerani and Spelt, 2010); describing these different relations is beyond the scope of the present paper.

### 1.3 Shape Evolution in the Chipping Regime

5 In the limit of pure chipping, where the fragments produced by each impact are suitably small compared to the parent clast, the shape evolution of a particle can be modeled purely geometrically. Essentially, areas that protrude from the pebble have a positive curvature and are more likely to strike another particle or the bed surface and chip off. Thus, particles undergo curvature-driven attrition that evolves their shape toward a sphere (Firey, 1974; Domokos et al., 2009; Várkonyi and Domokos, 2011; Domokos et al., 2014). In the typical case of bed load with gravels impacting a streambed, the situation is close to the  
10 purely curvature-driven limit (Szabó et al., 2013, 2015; Novák-Szabó et al., 2018). The particle shape evolution resulting from attrition by chipping occurs in two phases (Domokos et al., 2014). The first phase involves an angular clast quickly becoming round as protruding edges are removed without a significant change in axis dimensions (Krumbein, 1941; Kuenen, 1956; Adams, 1979; Domokos et al., 2014). The second phase occurs once a particle is entirely convex as the axis dimensions are slowly reduced (Domokos et al., 2014). Field studies indicate that two-phase attrition applies to sediment in both fluvial and  
15 aeolian environments (e.g., Szabó et al., 2013; Miller et al., 2014; Novák-Szabó et al., 2018). It has been found that circularity and aspect ratio are two convenient shape parameters that effectively characterize this shape evolution (Miller et al., 2014); we adopt them in this study.

In order to observe the shape changes in particles undergoing attrition on a spectrum from chipping to fragmentation, one can vary either material strength or impact energy (Eq. 2). In this study, concrete particles with the same initial mass were  
20 repeatedly dropped from the same height in a rotating drum to simulate transport. Thus, initial impact energy remained constant and experiments were conducted for varying material strength ( $A_b$ ). Mechanically strong particles experienced chipping with nearly constant mass loss per impact on their trajectory toward a spherical shape. Weak particles would fragment into large, irregular pieces after as few as 5 impacts. Concrete particles with intermediate strengths experienced both chipping and fatigue



failure, becoming more rounded but with occasional fragmentation events that prevented the particles from following the expected “universal rounding” curve associated with chipping.

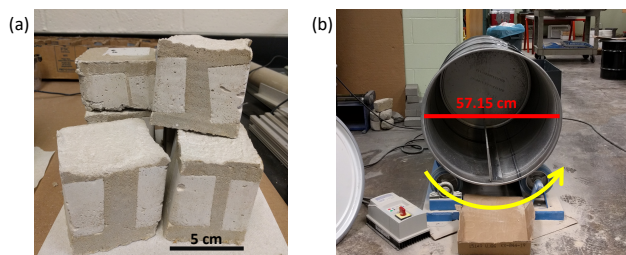
## 2 Methods

Experiments were conducted by placing concrete blocks of varying strength in a rotating drum. The concrete blocks were  
5 created by pouring a mixture of concrete mix and sand into 6-cm cubical molds (Fig. 3a). While concrete mixes meant for  
general construction purposes usually consist of gravel, sand, and Portland cement, the concrete mix used in the experiments  
was sieved to remove all clasts larger than coarse sand. The ratio of concrete mix to sand was varied in order to create particles  
of different strength with varying material susceptibility parameters ( $A_b$ ) that would undergo chipping, fragmentation, or a  
combination of the two breakdown mechanisms in the rotating drum. A total of 10 different mixtures were created—80 %, 75  
10 %, 66.7 %, 50 %, 33.3 %, 25 %, 20 %, 16.7 %, 14.3 %, and 12.5 %—where percentage indicates volume of concrete mix in  
the mixture (VCM). There were 5 particles created for each mixture, resulting in a total of 50 concrete particles tested.

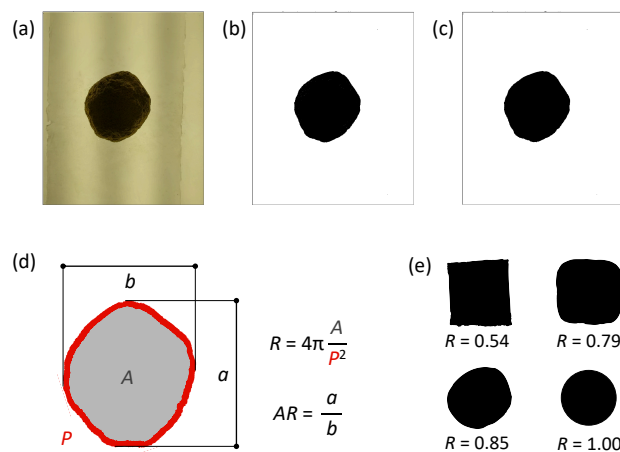
After the concrete particles had dried in their molds, they were then placed in a 208.2 L (55 gal.) open head carbon steel  
drum with a lid fastened by a bolt ring closure (Fig. 3b). The drum was 57.15 cm in diameter and 83.31 cm long. A steel paddle  
was attached to the interior of the drum to ensure that particles would collide with the side of the drum, as opposed to sliding  
15 along the bottom during the experiment. All experiments were run with the drum rotating in the counterclockwise direction at  
20 rotations per minute. Each particle was rotated in specific intervals according its composition. Weaker particles (12.5-20 %  
VCM) were removed after every rotation, intermediate particles (25-50 % VCM) were rotated in intervals from 5 to 20 drops,  
and stronger particles (66.7-80 % VCM) remained in the drum for intervals ranging from 50 to 500 drops. After each rotation  
interval, the particle was removed from the drum, weighed, and photographed before being returned to the drum to undergo  
20 another series of rotations.

At the conclusion of the rotating drum experiments, all images were analyzed using ImageJ. The image processing program  
converted the original photograph into a binary image in order to isolate the shape of the particle and measure shape parameters,  
including area, perimeter, circularity, and aspect ratio (Fig. 4). To verify the circularity measurements calculated by the image  
processing program, the shape measurement algorithm was applied to synthetic circles and squares of known shape. While  
25 measured circularity was found to be resolution dependent, the maximum error was 10% over the resolution range that is  
relevant for our experiments.

An Instron Universal Testing System was utilized to conduct uniaxial compression tests of the concrete particles. Similar  
to the rounding experiments, 10 mixtures of concrete mix and sand were used with 5 particles created from each mixture.  
A 25 kN load cell was used for the particles made of 12.5% to 50% VCM and a 150 kN load cell was used for particles  
30 with 66.7%, 75%, and 80% VCM. During the compression tests, the upper plate was driven down at a constant rate of 3 mm  
 $\text{min}^{-1}$  with maximum deformation set at 6 mm. Thus, the compression experiments produced a series of stress-strain curves  
to demonstrate the strength of the various concrete mixtures. Yield strength—the transition from (linear) elastic behavior to  
plastic deformation on a stress-strain curve—is frequently used as a measure of material strength. Our stress-strain curves,



**Figure 3.** Experimental materials and setup. (a) Concrete particles prior to breakdown in the drum. (b) The 208.20 L (55 gal.) carbon steel drum used for collision experiments. The drum is 57.15 cm in diameter, 83.31 cm long, and rotates counterclockwise at 20 rotations per minute.



**Figure 4.** A visual depiction of the ImageJ macro used to determine shape parameters. First, a photograph was taken of the particle at a fixed distance over an LED light table. The macro (a) cropped the image to include only the particle, (b) converted the original image into a binary image, then (c) filled any holes within the particle shape. The macro then (d) measured the area ( $A$ ), perimeter ( $P$ ), major axis ( $a$ ), and minor axis ( $b$ ) of the particle. These measurements were used to calculate circularity ( $R$ ) and aspect ratio ( $AR$ ), parameters that were used to quantify shape change over the course of the rounding experiments. Circularity measures how closely a shape approaches that of a circle, where  $R = 1$  indicates a perfect circle and  $R < 1$  indicates deviations from a circle. Part (e) shows circularity measurements for particles of different shapes.

however, were highly variable; especially for weaker materials that did not follow classic brittle failure (Fig. 5a). Accordingly, here we use ultimate strength—the greatest stress withstood by a material—similar to some previous studies (e.g., Sklar and Dietrich, 2001; Miller and Jerolmack, 2020). It is assumed that ultimate strength is proportional to yield strength. Young’s modulus is typically determined from a linear fit to the stress-strain plot — i.e., in the elastic regime before failure. Due to the variable shapes of our stress-strain curves, however, we estimated Young’s modulus as the ultimate strength divided by the associated strain (Fig. 5a), in order to avoid ambiguity of how to choose an approximately linear regime over which to fit.



### 3 Results

#### 3.1 Material Properties

An analysis of uniaxial compression test results revealed that particles with a higher percentage of concrete mix have a greater material strength and stiffness, as measured through ultimate strength and Young's Modulus (Fig. 5a). Mechanically strong particles could withstand loads ranging from  $3.0 \times 10^6$  to  $5.4 \times 10^6$  N m<sup>-2</sup>, although the strongest particle (66.7% VCM) held up to  $1.2 \times 10^7$  N m<sup>-2</sup>. Mechanically weaker particles could hold loads ranging from  $1.0 \times 10^5$  to  $1.2 \times 10^6$  N m<sup>-2</sup>, although the weakest material (12.5% VCM) withstood as little as  $4.3 \times 10^4$  N m<sup>-2</sup> before failing (Fig. 5b). Additionally, there was considerable variation in the measured material properties within samples of a specified composition. Despite this variation, the general trend indicates a positive relation between ultimate strength and proportion of cement mix.

A similar trend applies to the particles' resistance to deformation, as measured through Young's Modulus. Mechanically strong particles had Young's Moduli ranging from  $8.4 \times 10^7$  to  $1.6 \times 10^8$  N m<sup>-2</sup>, although the strongest particle (66.7% VCM) had a Young's Modulus measured at  $3.8 \times 10^8$  N m<sup>-2</sup>. Mechanically weaker particles had Young's Moduli ranging from  $1.5 \times 10^6$  to  $1.1 \times 10^7$  N m<sup>-2</sup>, although the weakest material (12.5% VCM) had a Young's Modulus measured at  $2.7 \times 10^5$  N m<sup>-2</sup> (Fig. 5c). Despite variation in the measured Young's Moduli, the general trend indicates that materials with a higher percentage of concrete mix were more resistant to deformation.

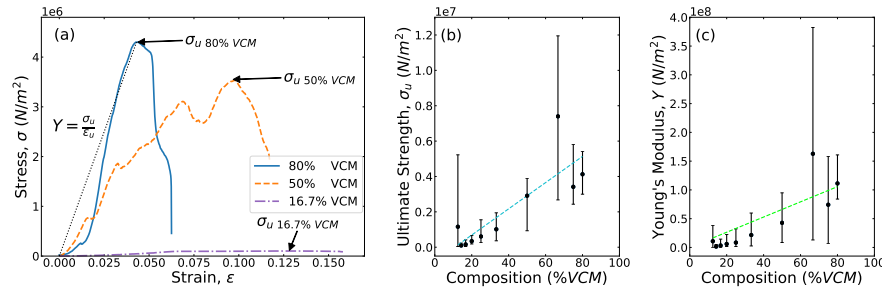
The density of each type of concrete particle was also calculated from measured weight, and volume for each cube estimated from axis dimensions of the molds. Density did not vary systematically with material strength, however; the average density was approximately  $\rho_s = 2000$  kg/m<sup>3</sup>, with a range of  $\pm 20\%$ . Given the error bars in the measurements of ultimate strength and Young's Modulus, we use the best-fit lines from Figure 5a and 5b to determine the associated material properties for each composition. This is done to minimize multiplicative errors in calculating the brittle Attrition Number,  $A_b$ .

#### 3.2 Mass Loss

The mass of each particle decreased as that particle rotated in the metal drum. To quantify and characterize mass loss, we utilize the following parameters: mass fraction ( $M$ ), cumulative mass loss ( $\mu$ ), and fractional mass loss per impact ( $\Delta m_*$ ). Mass fraction is defined as the ratio of the mass of the particle during a given rotation to the initial particle mass ( $M = m_i/m_o$ ). Cumulative mass loss is the ratio of remaining mass to the initial particle mass ( $\mu = 1 - M$ ). Fractional mass loss per impact is the ratio of the mass lost during a given impact to the particle mass just prior to that impact ( $\Delta m_* = (m_{i-1} - m_i)/m_{i-1} = \Delta m/m_{i-1}$ ).

As particles experienced transport in the metal drum, mass was reduced most rapidly at the beginning of each experiment (Fig. 6a, c, e). For mechanically strong particles, fractional mass loss per impact was approximately constant (Fig. 6b, d, f). The fractional mass loss per impact observed in intermediate particles and weak particles was variable and fluctuated over the duration of the experiments. The weakest particles sustained the greatest fluctuations in mass loss, indicating that these particles experienced multiple fragmentation events as they rotated in the drum.

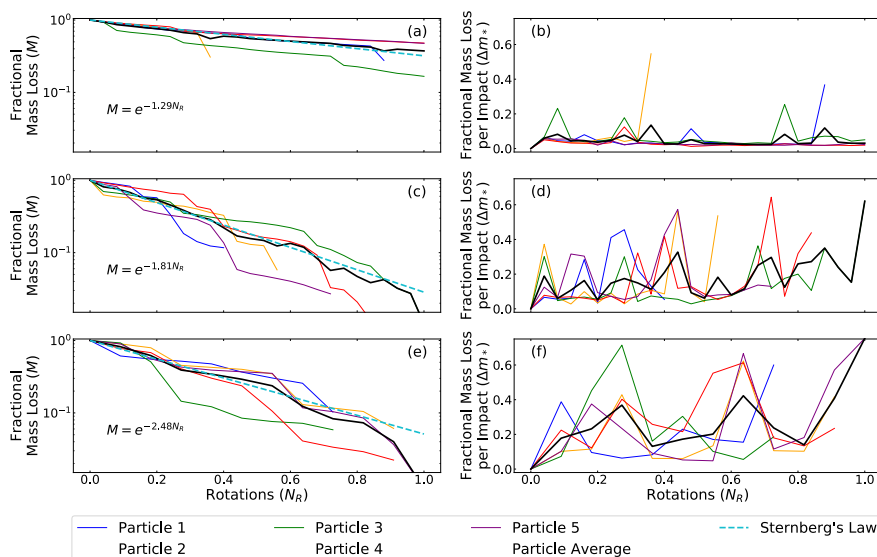




**Figure 5.** Material strength properties determined by uniaxial compression tests. (a) Stress-strain curves for concrete particles with varying proportions of concrete mix. Ultimate strength,  $\sigma_u$ , is defined as the greatest stress withstood by the particle and is indicated on the stress-strain figure. Young's Modulus,  $Y$ , is calculated as the ultimate strength divided by the corresponding strain ( $Y = \sigma_u / \epsilon$ ). (b) Average ultimate strength ( $n = 5$ ) for each particle composition, where error bars indicate the range of values produced by experiments. A line of best fit was used to determine the material strength of each particle composition used in the study and can be described through the following equation:  $\sigma_u = 7.41e4C_m - 7.94e5$ , where  $C_m$  is the composition of the concrete particle based on the volume concrete mix (VCM);  $R^2 = 0.69$ . (c) Average Young's Modulus ( $n = 5$ ) for each particle composition, where error bars indicate the range of values produced by experiments. A line of best fit was used to determine the Young's Modulus for each particle composition and can be described through the following equation:  $Y = 1.32e6C_m$ , where  $R^2 = 0.67$

In order to visualize mass loss distribution, a series of histograms were generated for each particle strength (Fig. 7). These histograms indicate that mechanically strong particles tended to produce chips with a narrow distribution of masses that are small in comparison to the mass of the initial particle. Mechanically weak particles tended to produce fragments with a wide distribution of masses that could be a significant proportion of the initial particle mass.

5 We now test the control of material strength on mass loss across the chipping to fragmentation transition. We find first that our data do not follow the inverse-square relation between mass loss and ultimate strength proposed by Sklar and Dietrich (2001, 2004). Instead, we find an exponent of less than (but close to) one, indicating that ultimate strength alone is insufficient to describe the control of material on attrition rate (Fig. 8). To make contact with the proposed relations Eqs. 1 and 2, we first fit an exponential Sternberg-like relation to the curves of fractional mass loss against rotation number, and verify that the data are reasonably well fit (Fig. 6). This supports the notion that mass lost per unit energy is constant (Eq. 2) — even if highly variable for weaker materials. We then examine Eq. 2 and rearrange to solve for  $k = AC_1 = 2\Delta m_*/v_i^2$  to empirically determine  $k$  for each material. We anticipate that the empirical mass-loss parameter  $k$  scales linearly with the brittle Attrition number  $A_b$ ; indeed the data are consistent with such a relation (Fig. 8b) and with recent experiments on binary collisions of real rocks Miller and Jerolmack (2020), though the experimental parameter  $C_1$  determined here is three orders of magnitude larger than  
 15 that study. A secondary pattern is that, for the strongest materials, mass loss is independent of material strength (Figs. 7, 8).

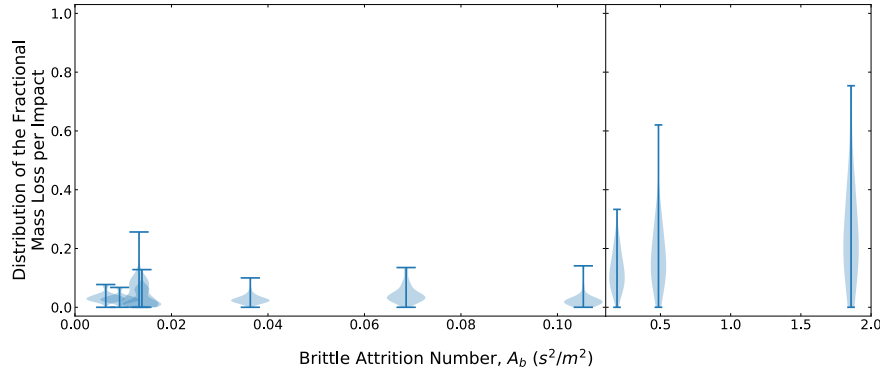


**Figure 6.** Mass loss parameters over the course of rotation experiments. Figures to the left show changes in mass fraction,  $M$ , for (a) 80 % VCM, (c) 50 % VCM, and (e) 16.7 % VCM. Figures to the right show fractional mass loss per impact,  $\Delta m_*$ , for the same materials in (b), (d), and (f). All 5 particles of a given composition are represented on each plot, with the average ( $n = 5$ ) depicted as a thicker black line. The number of rotations undergone by each particle have been normalized by the total number of rotations to allow comparison—particles with 80 % VCM underwent 4400 rotations, particles with 50 % VCM underwent 3500 rotations, and particles with 16.7 % VCM underwent 25 rotations. Mechanically weak particles, represented at the bottom of the figure, lost significant amounts of mass during each rotation and continued to be placed in the drum until very little of the original particle remained. Additionally, Sternberg’s Law (Eq. 1)—denoted by the cyan dashed line—was fit to each plot of fractional mass loss with the form  $M = m_i/m_o = e^{-k_{C_m} N_R}$ , where  $k_{C_m}$  is a material parameter that differs for each concrete mixture. We can see that fractional mass loss scales exponentially with the number of impacts in the rotating drum.

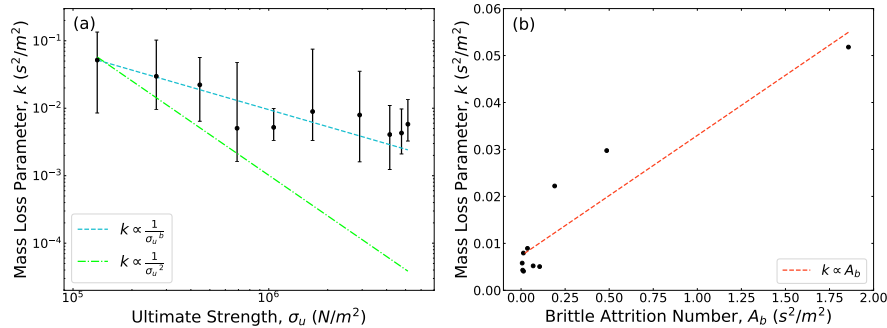
### 3.3 Particle Shape

Rounding experiments showed that particles of differing mechanical properties experienced different shape evolution trajectories. Mechanically strong particles with a greater proportion of concrete mix evolved from a cubical block toward a sphere. The mechanically weak particles began with the same cube-like shape, but fragmented into several angular pieces during their  
 5 time in the rotating drum (Fig. 9). The strong particles could withstand thousands of impacts (66.7–80 % VCM), while the weakest particles experienced 10–20 impacts before disintegrating (12.5–16.7 % VCM). Particles of intermediate strength (20–50 % VCM) evolved from a cubical block toward a spherical shape, but experienced several large breakage events along the trajectory toward a sphere. These intermediate particles remained in the rotating drum for several hundred impacts.

Circularity, a measure of how closely a shape approaches a circle, increased over the course of the experiments. The cir-  
 10 cularity of mechanically strong particles increased smoothly and monotonically over time toward a spherical shape ( $R = 1$ ).

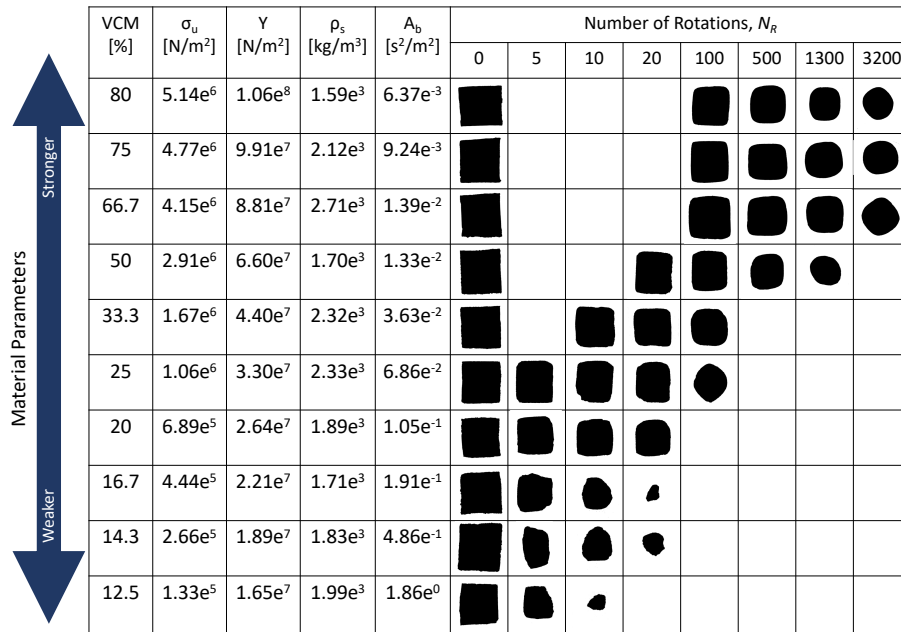


**Figure 7.** Distribution of the fractional mass loss per impact,  $\Delta m_*$ , for all material strengths, where the distribution is an average for all particles of a given composition ( $n = 5$ ). Mechanically weak particles (higher  $A_b$ ) produced fragments in a wide range of sizes and weights, while mechanically strong particles (lower  $A_b$ ) generated a narrow distribution of chip sizes. This transition occurs around  $A_b = 0.11 \text{ s}^2\text{m}^{-2}$ . Additionally, there is a second transition for brittle Attrition Numbers lower than  $\sim 0.01 \text{ s}^2\text{m}^{-2}$ , where the size distribution of fragments narrows further.



**Figure 8.** (a) The average mass loss parameter ( $n = 5$ ) for all material strengths, where  $k = AC_1 = 2\Delta m_*/v_i^2$ . The cyan dashed line indicates an inverse power law line of best fit ( $k = 1.03e3/\sigma_u^{8.38e-1}$ , where  $R^2 = 0.94$ ). The green dashed line indicates an inverse square law ( $k = 1.02e9/\sigma_u^2$ , where  $R^2 = 0.64$ ). The inverse square relationship was proposed by Sklar and Dietrich (2004) to describe brittle fracture. Error bars indicate the range of values produced by experiments. (b) Relation between the empirically-determined mass loss parameter,  $k$ , and the brittle Attrition Number determined from material properties,  $A_b$  (Miller and Jerolmack, 2020); the best-fit linear relation is  $k = 0.026A_b + 0.0074$ , where  $R^2 = 0.88$ . Thus experimental constant  $C_1 = 39.03$ .

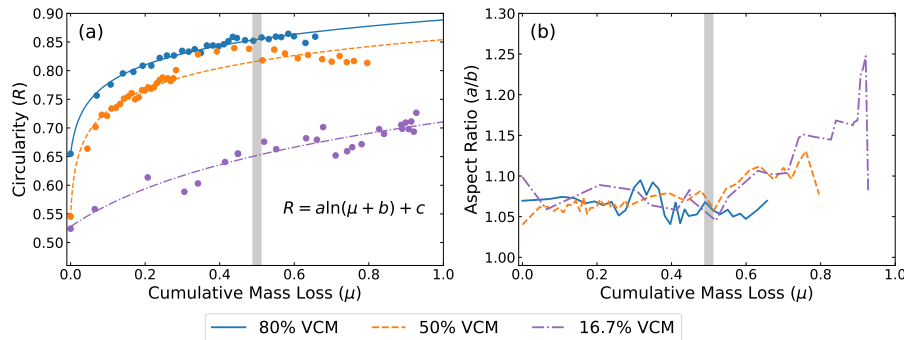
For intermediate and weak particles, circularity increased, but exhibited significant fluctuations associated with fragmentation events. Mechanically strong particles reached circularity values around 0.85, while weak particles reached values around 0.65 (Fig. 10a). Visual inspection confirms these trends (Fig. 9). The circularity associated with different particle strengths was examined across experiments by utilizing circularity values measured at 50 % cumulative mass loss. This comparison indicates



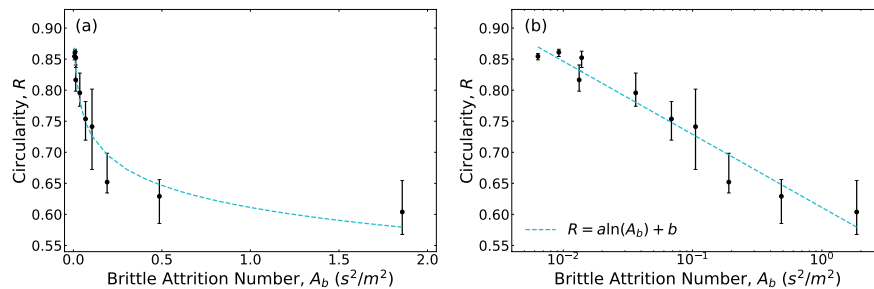
**Figure 9.** Progression of particle breakdown according to material parameters and approximate number of impacts. The material parameters of each particle type are indicated using volume concrete mix (VCM), ultimate strength ( $\sigma_u$ ), Young’s Modulus ( $Y$ ), and brittle Attrition Number ( $A_b$ ). Also shown is particle density ( $\rho_s$ ), which does not vary systematically with material strength. Mechanically strong particles achieved highly rounded shapes over thousands of rotations, while mechanically weak particles quickly fragmented into irregular shapes over 10–20 impacts. Intermediate particles exhibited both behaviors.

that there is a roughly logarithmic relation between circularity and brittle Attrition Number (Fig. 11). Data also suggest that a limit exists where particles achieve similar circularity values regardless of material strength. This observation aligns with geometric theory, which predicts that all particles undergoing chipping will reach the same circularity value for a given amount of mass lost (Novák-Szabó et al., 2018).

- Aspect ratio, a ratio of the major and minor axes, was also used to visualize particle shape change over the course of the experiments. An aspect ratio of 1 indicates that the length and width of a particle are equal. Since all particles were formed in cubical molds, the aspect ratio of the particles before the tumbling experiments was close to 1. Since the mechanically strong particles evolved toward a sphere, the aspect ratio remained close to 1 for the duration of the experiment. The aspect ratio of intermediate and weak particles deviated from 1 over the course of the experiments, indicating that fragmentation events were occurring (Fig. 10b).



**Figure 10.** Evolution of shape parameters as particles undergo transport in the rotating drum. As particles are rotated, they continually lose mass and become rounder over time. (a) Average circularity ( $n = 5$ ) for all particles increases during transport, with mechanically strong particles achieving higher circularity values. The relationship between circularity and cumulative mass loss can be explained using a logarithmic relationship ( $R = a \ln(\mu + b) + c$ ). (b) Average aspect ratio ( $n = 5$ ) begins close to 1 for all particle strengths and remains low for mechanically strong particles. However, aspect ratio can fluctuate wildly as mechanically weak particles undergo transport. The grey region in both figures indicates 50 % cumulative mass loss and the corresponding circularity and aspect ratio. Parameters at 50 % cumulative mass loss are used to compare particle characteristics across experiments.



**Figure 11.** Average circularity ( $n = 5$ ) at 50 % cumulative mass loss for particles of different mechanical strengths in (a) linear space and (b) semi-log space. The logarithmic relationship between circularity and brittle Attrition Number can be described through the following equation:  $R = -0.051 \ln(A_b) + 0.61$ , where  $R^2 = 0.95$ . Error bars indicate the range of values produced by experiments.

#### 4 Discussion

Tests indicate that material strength and resistance to deformation increased with the amount of concrete mix in the particle's composition. However, there was significant variation in the ultimate strengths and Young's Moduli recorded for each set of particles, as well as the shapes of the stress-strain curves (Fig. 5a). Typical stress-strain curves show an elastic region, where the curve follows a linear regime, and a plastic region, where the curve rounds and eventually reaches the fracture point. Some of our curves do not follow this pattern, but instead display a series of peaks in the transition from the elastic to the plastic regime (Fig. 5a). This variation in particle strength and behavior, in response to compressional stress, is likely due to inconsistencies

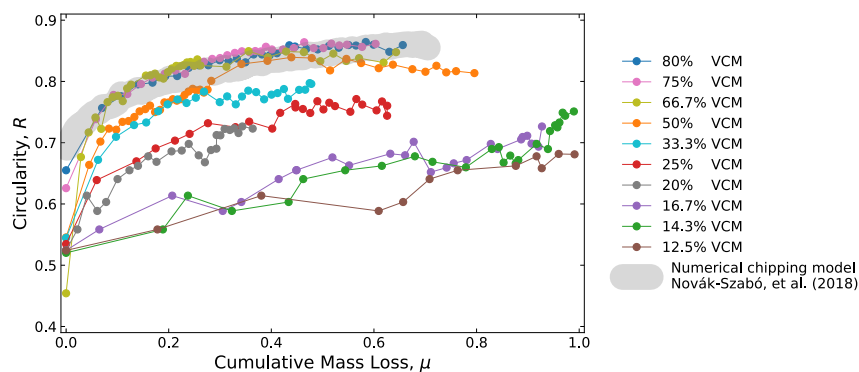


inherent in the process of making these concrete particles. It is possible that inhomogeneities produced during particle creation make the curves deviate from industrially produced concrete (e.g., Lan et al., 2010). Furthermore, molds used to create the concrete particles varied in shape from 6 cm to 8 cm, and asperities would form on the concrete along the open side of the mold. These asperities may also have influenced compression testing by reducing the ultimate strength of each particle. The  
5 compressing plates should come into contact with flat surfaces; asperities would cause the force to load unevenly and may result in premature failure and the observed peaks in the stress-strain curves (e.g., Vasconcelos and Lourenço, 2009).

In addition to potentially causing premature failure or unusual patterns of strain in particles, the variation in initial particle shape may have impacted the rounding experiments. Ideally, particle shape should not impact the outcome of rounding experiments because particle shape does not play a key role in setting attrition rate (Várkonyi and Domokos, 2011; Szabó et al.,  
10 2013; Bertoni et al., 2016). Particle composition, on the other hand, can influence shape evolution (Sklar and Dietrich, 2004). The high composition of sand in mechanically weak particles frequently caused fragments to disintegrate on impact, producing a large population of sand and fine concrete fragments that may contribute to the bimodal distribution seen in the mass loss distributions of some intermediate strength particles (Fig. 7). The population of fine fragments was not collected or measured, as this pattern of disintegration differs from fragmentation in natural materials.

15 Nonetheless, we can identify a clear control of material properties on the attrition rate, that is consistent with the brittle Attrition Number,  $A_b$ , being the relevant material grouping. The findings are in agreement with recent results on binary collisions of natural rock materials Miller and Jerolmack (2020). The large difference of the experimental parameter  $C_1$  between that study and ours is not entirely understood. This parameter likely encodes details of the collision process (impact angle, rotation, etc.) but also, by definition, includes details of how material properties were measured. Our findings reinforce the idea that the  
20 brittle Attrition Number is a useful material parameter for determining relative susceptibility to attrition, but uncertainties in  $C_1$  limit our ability to directly extrapolate experimental findings to the field. Additionally, results indicate that there is an upper limit where material strength no longer influences particle breakdown. Our data indicate that the brittle Attrition Numbers below  $0.014 \text{ s}^2\text{m}^{-2}$  are approaching this limit. Both circularity and mass loss become nearly constant for particles with brittle Attrition Numbers lower than the limit (Fig. 8a, 11).

25 The results of the rounding and compression experiments confirm the existence of a pure chipping regime. Other studies have utilized both experiments and field observations from rivers, beaches, and dune fields to predict the shape evolution of particles undergoing attrition by chipping (e.g., Szabó et al., 2013; Miller et al., 2014; Bertoni et al., 2016; Novák-Szabó et al., 2018). The circularity of mechanically strong particles produced by this study increases smoothly and monotonically with increasing mass loss. Additionally, circularity as a function of cumulative mass loss for our relatively strong particles plot very  
30 close to the universal curve found from experiments and field observations on natural materials taken in diverse environments (Fig. 12). Deviations from this universal curve in the initial stage of mass loss are likely due to the different initial conditions; natural rock fragments have common initial shapes (Domokos et al., 2015, 2020), while the particles produced in this study are initially cubes with variable asperities on their faces. Additionally, particles with  $\text{VCM} < 50 \%$  begin to show significant deviations from the universal curve; circularity grows slowly and more erratically. We view this departure, which increases as  
35 the material strength weakens, as a direct signature of the prevalence of attrition by fragmentation. Particles with  $\text{VCM} > 50$



**Figure 12.** Circularity as a function of cumulative mass loss for all material strengths tested, where circularity measurements are averaged for all particles of a given composition ( $n = 5$ ). Mechanically strong particles (66.7–80 % VCM) conform to the universal curve described by Novák-Szabó et al. (2018), which is identified by the grey region.

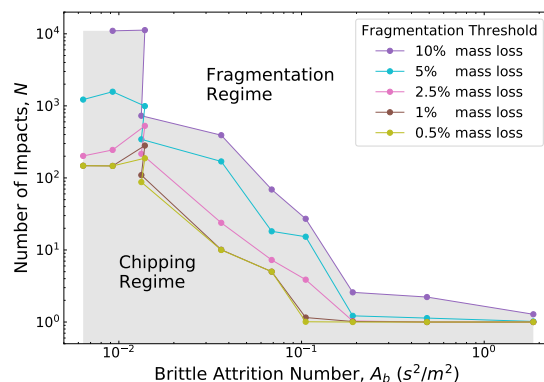
% conform to the universal curve for circularity as a function of mass loss, suggesting that materials with a brittle Attrition Number of  $0.014 \text{ s}^2\text{m}^{-2}$  or lower are in the limiting regime of pure chipping.

For the attrition of materials, the two end-members can readily be defined. In the limit of pure chipping, repeated low-energy impacts result in small-scale shattering that is confined to the near surface of the impacted particle (Miller and Jerolmack, 2020).

- 5 This manifests in the patterns of rounding observed in river rocks transported by bed load. In the limit of pure fragmentation (catastrophic failure), where a single high-energy impact is capable of propagating cracks throughout the body of a rock, bulk breakup produces large angular fragments from the initial particle. The intermediate regime, however, is more difficult to determine. Here, fatigue failure occurs as multiple impacts are required to propagate cracks through the bulk of the particle. By performing experiments that span these regimes, we can populate the theoretical phase-space for attrition conceptualized by
- 10 Zhang and Ghadiri (2002) (Fig. 1). By setting a threshold for fragmentation based on fractional mass loss per impact, we can then approximate the number of impacts required for a material to fragment (Fig. 13). In defining this threshold as a significant fractional mass loss per impact (e.g., 10 %) (Shipway and Hutchings, 1993), then only the weakest material (12.5 % VCM) that breaks with nearly every impact can be categorized in the catastrophic failure regime. Alternately, if we define the threshold between chipping and fragmentation as a low fraction of the mass lost (e.g., 1 % or 0.5 %), then the four weakest concrete
- 15 mixes (12.5 %, 14.3 %, 16.7 %, and 20 % VCM) fall within the catastrophic failure limit.

## 5 Conclusions

This study connected attrition mechanism to the resulting shape evolution in order to better understand the transition from chipping to fragmentation. By simulating transport over a wide range of material strengths, we were able to populate a phase space for attrition. Concrete particles were rotated in a metal drum to simulate transport, and results indicate that mechanically



**Figure 13.** Approximate number of impacts required to fragment particles of different material strength. The fragmentation threshold, or boundary between the chipping and fragmentation regimes, is defined by the fraction of mass lost during a given impact (e.g., fragmentation occurs if 10 % of a particle’s current mass is lost).

strong particles evolve smoothly and monotonically toward a spherical shape, while weak particles rapidly break into irregular, angular pieces. Intermediate particles undergo a combination of the two breakdown mechanisms and erratically evolve toward a rounded shape. Our finding that the brittle Attrition number is the relevant material grouping governing attrition rates supports recent experimental results from binary collisions of natural rock materials (Miller and Jerolmack, 2020), and indicates that  $A_b$  is a useful similarity parameter for scaling experiments.

Analysis of shape and mass loss parameters indicate that a limit exists where material strength has limited influence on particle breakdown. In the limit determined by this study—a brittle Attrition Number of  $0.014 \text{ s}^2\text{m}^{-2}$  or lower—circularity and fractional mass loss approach constant values. This is the pure chipping limit of impact attrition. This condition appears to be common, even universal, for natural rocks undergoing bed-load transport (Novák-Szabó et al., 2018). The agreement of our experiments using concrete blocks, with previously described universal rounding, affirms the robustness of geometric shape evolution by chipping and also the relevance of our experiments to natural pebbles. As material strength decreased in our experiments, the shape evolution of each particle began to deviate from the universal shape evolution curve, with intermediate and weak particles experiencing a greater variation in mass loss per collision and reaching lower circularities over the course of transport in the metal drum. This phase space delineates the “fatigue failure” regime of pebbles, where multiple collisions are required to grow cracks in the bulk to the point where they span the sample and induce fragmentation. While fatigue failure is often invoked to model bedrock erosion due to impact attrition (Sklar and Dietrich, 2001, 2004), it has not to our knowledge been delineated in this manner before. This fatigue failure regime can be thought of as representing a continuous phase transition from pure chipping to pure fragmentation, where the probability of fragmentation per collision increases as material strength weakens (for a constant impact energy). This picture is consistent with the probabilistic conceptual model of Zhang and Ghadiri (2002), where this transition was cast as a function of impact energy (for constant material strength).





Sufficiently weak materials (or large collision energies) will catastrophically fail with a single impact; this limit is likely reached in nature for rockfalls and landslides. This work suggests that the parameter grouping  $A_b \Delta E$  may serve to predict the attrition regime.

*Data availability.* All relevant data used to generate the figures in this paper are deposited in figshare, a free and open repository. These  
5 datasets can be accessed using the following link: <https://doi.org/10.6084/m9.figshare.14069105>.

*Author contributions.* S.B. conducted the experiments, led data analysis, and wrote the paper. D.J.J. supervised the research, assisted in analysis and interpretation of data, and edited the paper.

*Competing interests.* The authors declare that they have no conflict of interest.

*Disclaimer.* The authors declare no conflicts of interest.

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